

A counter-top search for macroscopic dark matter

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(Dated: October 8, 2019)

A number of dark matter candidates have been discussed that are macroscopic, of approximately nuclear density, and scatter ordinary matter essentially elastically with approximately their geometric cross-section. A wide range of mass and geometric cross-section is still unprobed for these “macros.” Macros passing through rock would melt the material in cylinders surrounding their long nearly straight trajectories. Once cooled, the resolidified rock would be easily distinguishable from its surroundings. We discuss how, by visually examining ordinary slabs of rock such as are widely available commercially for kitchen countertops, one could probe an interesting segment of the open macro parameter space.

There is considerable, and widely known, evidence for the existence of dark matter [1] or for the need to modify General Relativity (see [2] for a pedagogical review of modified gravity and [3] for a more recent review on the subject). Macroscopic dark matter (macros) represents a wide class of alternatives to particle dark matter – large objects, probably composites of fundamental particles, that are “dark” because their large mass implies a low number density and a small geometric cross-section per unit mass, even though the cross-section of each object is large. There remains a large range of macro mass M_X and geometric cross-section σ_X that is still unprobed by experiments or observations.

A most intriguing possibility is that macros are made of Standard Model quarks or baryons bound by Standard Model forces. This suggestion was originally made by Witten [4], in the context of a nuclear bag model and a then-possible first-order QCD phase transition. A more realistic version was advanced by Lynn, Nelson and Tetradis [5] and Lynn again [6], who argued in the context of $SU(3)$ chiral perturbation theory that “a bound state of baryons with a well-defined surface may conceivably form in the presence of kaon condensation.” Nelson [7] studied the possible formation of such “nuggets of strange baryon matter” in an early-universe transition from a kaon-condensate phase of QCD to the ordinary phase. Others have suggested non-Standard Model versions of such nuclear objects and their formation, for example incorporating the axion [8].

Such objects would presumably have densities that are comparable to nuclear density (which we take to be $\rho_{nuclear} = 3.6 \times 10^{14} \text{ g cm}^{-3}$). This is much higher than ordinary “atomic density” ($\rho_{atomic} = 1 \text{ g cm}^{-3}$), and much lower density than black holes. The unconstrained macro parameter space includes a wide range of M_X for macros of nuclear density (see [9] and [10]). For $M_X \leq 55 \text{ g}$, careful examination of specimens of old mica for tracks made by passing dark matter [11, 12] have ruled out such objects as the primary dark matter candidate (see Figure 1). For $M_X \geq 10^{24} \text{ g}$, a variety of

microlensing searches have similarly constrained macros [13–16]. For $M_X \gtrsim 10^{15} \text{ g}$, macros incident on white dwarfs would trigger thermonuclear runaways [17], as previously shown for primordial black holes [18], and are ruled out. Reference [19] utilized the full Boltzmann formalism to obtain constraints from macro-photon elastic scattering using the first-year release of cosmic microwave background data from the Planck satellite. Although strictly we should differentiate between the macro-baryon cross-section and the macro-photon interaction, these are macroscopic cross-sections and it is difficult to imagine that these will differ significantly unless the photon wavelength is at least comparable to the size of the macro. We therefore take the constraint in [19] on $\sigma_{macro-photon}$ to apply to σ_x directly, for $\sqrt{\frac{\sigma_x}{\pi}} \geq \lambda_{CMB} \sim 1000 \text{ nm}$ (for photons at recombination). For $\sqrt{\frac{\sigma_x}{\pi}} \leq \lambda_{CMB}$, wave optics or quantum-mechanical effects may weaken the constraints in a model-dependent way. We hatch the corresponding region of parameter space to highlight this effect. Prior work had already constrained a similar range of parameter space by showing that the consequence of dark-matter interactions with Standard Model particles is to dampen the primordial matter fluctuations and suppress all structures below a given scale (see e.g. [20]). In between, $55 \text{ g} \leq M_X \lesssim 10^{15} \text{ g}$, the coast is so far clear for nuclear-density macro dark matter.

Recently, together with collaborators, we suggested how ultra-high-energy cosmic ray detectors that exploit atmospheric fluorescence could potentially be modified to probe parts of macro parameter space [21], including macros of nuclear density and intermediate mass. In this manuscript, we suggest how the approach applied to mica [11, 12] could be applied to a larger, widely available sample of appropriate rock, and used to search for larger-mass macros. (“Paleo-detectors” have also been considered recently for WIMPS [22, 23], and earlier [24] using ancient mica.) We also discuss our planned efforts to look for these tracks directly and the range of macro mass and cross-section that could be probed. For these purposes, we consider macros of a single mass and size,

even though a broad mass distribution is the expectation in the context of a composite dark matter candidate.

The energy deposited by a transiting macro through its elastic scattering off the rock is

$$\frac{dE}{dx} = \sigma_X \rho v_X^2, \quad (1)$$

where ρ is the density of the medium, σ_X is the geometric cross section of the macro and v_X is the macro speed.

For definiteness, we assume macros possess a Maxwellian velocity distribution

$$f_{MB}(v_X) = \left(\frac{1}{\pi v_{vir}^2}\right)^{\frac{3}{2}} 4\pi v_X^2 e^{-\left(\frac{v_X}{v_{vir}}\right)^2}, \quad (2)$$

where $v_{vir} \approx 250 \text{ km s}^{-1}$. This distribution is slightly modified by the Earth's motion. (See footnote on page 15 of [21] for more details.) The cumulative velocity distribution function is obtained by integrating f_{MB} up to the desired value of v_X . This allows us to determine the maximum mass M_X we can probe as a function of v_X .

The speed of a macro traveling through a medium is expected to evolve as

$$v = v_{X,0} e^{-\langle \rho x \rangle \frac{\sigma_X}{M_X}}, \quad (3)$$

where $\langle \rho x \rangle$ is the encountered column density. This will determine the maximum value of $\frac{\sigma_X}{M_X}$ expected to deposit sufficient energy to produce an observable signal.

If a macro were to pass through rock, the region nearest the trajectory would be ionized, the surrounding region vaporized, and an even larger region would be melted. The ionization and vaporization would result in extreme pressures, especially near the trajectory [25]. After resolidifying and cooling, the resulting rock would typically be petrologically distinguishable from the original rock around it. For example, even much lower energy-density lightning strikes on sand rich in silica or quartz are known to form fulgurites[26], glass tubes or clumps embedded in the sand. If the macro were large enough, the metamorphosed rock might even be visually distinguishable. For example, when light-colored granite is melted, it cools to form a dark obsidian-like stone. [27]

We solve the heat equation to find out to what distance a macro melts the rock surrounding its trajectory. This melt zone will cool to form a robust fossil record of the macro's passage. Unlike fulgurites, the track of a macro would be straight, and, unless the properties of the rock are highly anisotropic, the trail is likely to consist of a highly distinctive, long straight cylinder of circular cross section. This presents, in theory, a straightforward way of looking for macros in layers of rock.

Following the work of Cyncynates *et al.* [25], we approximate the initial heat deposition as a delta-function source along a straight line through the rock, and propagate the heat outwards according to the heat equation. The resulting time-dependent temperature field

$$T_c(r, t) = \frac{\sigma_X v_X^2}{4\pi\alpha c_p} \frac{e^{-\frac{r^2}{4t\alpha}}}{t}, \quad (4)$$

where σ_X is the macro cross section, v_X is the macro speed, c_p is the specific heat capacity of the rock, which we take to be $c_p \sim 900 \text{ J kg}^{-1} \text{ K}^{-1}$ [28], and α is its thermal diffusivity.

We invert (4) to obtain $\pi r(t, T_c)^2$, the area that gets heated to T_c (we take $T_c \approx 1200^\circ\text{C}$ [29]), such as the melting or vaporization temperature of the rock,

$$\pi r(t, T_c)^2 = 4\pi t \alpha \ln\left(\frac{\sigma_X v_X^2}{4\pi\alpha c_p T_c t}\right). \quad (5)$$

The maximum area that gets melted or vaporized is then

$$\pi r^2|_{max} = 4\pi t_{max} \alpha = \frac{\sigma_X v_X^2}{c_p T_c e}, \quad (6)$$

which occurs at a time

$$t_{max} = \frac{\sigma_X v_X^2}{4\pi\alpha c_p T_c e} \quad (7)$$

after the macro passage.

Note that this calculation ignores the specific heat of melting or vaporization, but we have checked that this is a small fraction of the deposited energy.

The expected number of macro passages through the rock depends on M_X .

$$N_{events} = \frac{\rho_{DM} A_{det} T_e v_X}{M_X} \approx 20 \left(\frac{A_{det}}{10m^2}\right) \left(\frac{T_e}{500My}\right) \left(\frac{kg}{M_X}\right), \quad (8)$$

where $\rho_{DM} \approx 5.3 \times 10^{-25} \text{ g cm}^{-3}$ [30], A_{det} is the cross-sectional area of the rock slab(s), and T_e is the exposure time, i.e. time the rock has spent near enough the surface to be exposed to the high-velocity macro flux.

In Figure 1, we present the regions of parameter space that could be probed by three possible searches, to highlight the different approaches that could be taken. We imagine a process that takes thin slabs of rock and inspects their surfaces for elliptical ‘‘melt patches’’ – the cross-sections of the circular melt cylinder. A macro would induce identical, aligned, elliptical patches on the front and back sides of a thin slab of uniform composition. The hypothesis that a dark elliptical region in a slab of light granite was caused by a macro could thus be substantiated by examining the obverse surface for a dark patch of matching size, ellipticity and alignment. There is no known alternate natural cause for such a pair of patches in granite.

Slabs of light-colored granite, such as are commonly used for kitchen counter tops, would seem to be ideal targets for inspection. These could be examined at commercial showrooms after they were polished, but before they were cut to size and installed. The uncut slabs are typically 2–3cm thick, and several square meters in area.

Melted regions of such granite will tend to be much darker and thus easy to see in the polished surface. They could then be confirmed by examining the back surface.

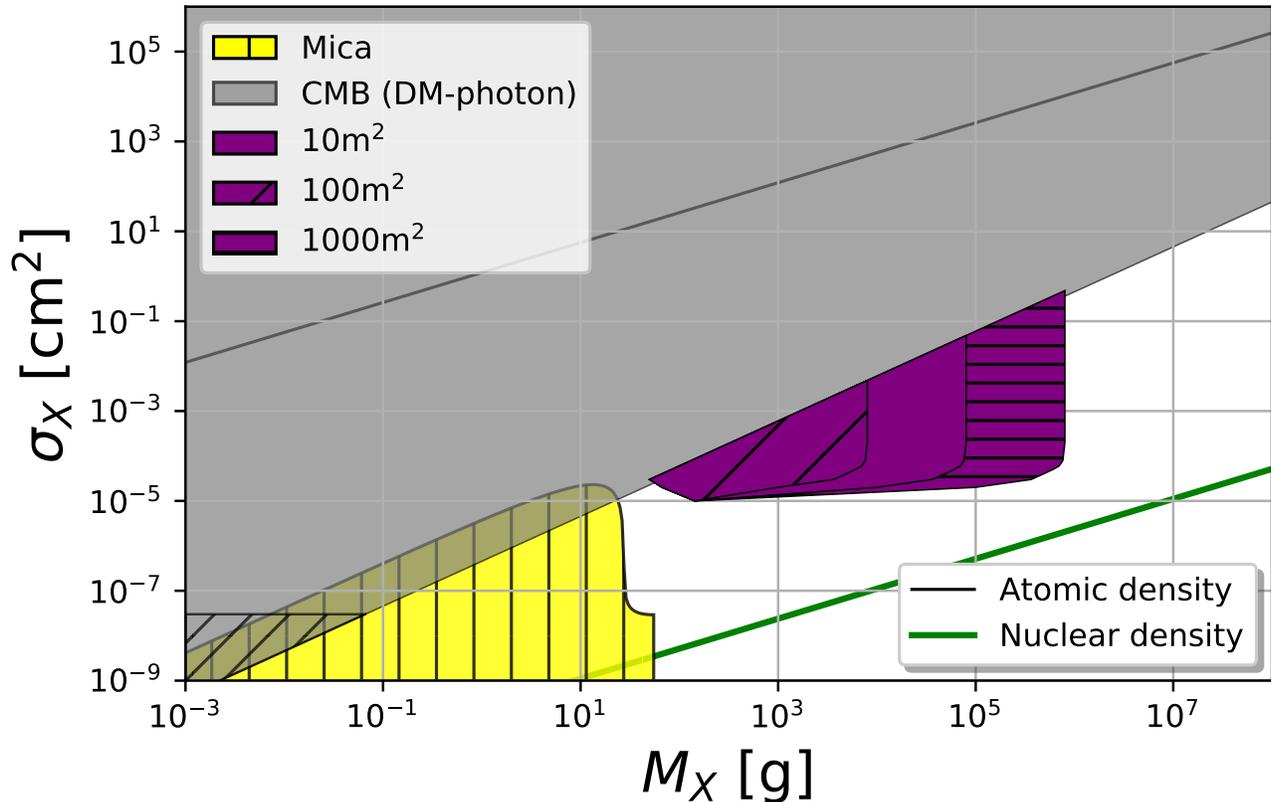


Figure 1. Illustrative regions of parameter space that could be probed from observations of slabs of ordinary rock are shown in purple. All three regions assume a minimum feature diameter of 1 mm. The region with diagonal hatching assumes a total slab area of 10m^2 , the unatched region is for a slab of area of 100m^2 and the horizontal-hatched region assumes a slab area of 1000m^2 . The region currently excluded by examination of ancient mica [11, 12] is shown in yellow with vertical hatching. The grey region is excluded from the effects of cosmic microwave background photons scattering off the macros [19]. The diagonal hatching of this bound in the bottom-left of the parameter space represents the region where the bound may weaken. Lines corresponding to nuclear and atomic density are shown for illustrative purposes.

As the cross-sectional area of the macro-induced melt-tube is proportional to the geometric cross-section of the macro, σ_X , the features caused by smaller macros are presumably harder to identify. On the other hand, lower-mass macros are more abundant. For a fixed (say nuclear) density this presents an obvious trade-off between feature size and feature abundance. The grain sizes in granite are often at the 1 mm scale or below. The smaller the grains in a given slab, the smaller the minimum feature size that can confidently be searched for in that slab. The rate of false detections can be determined from the rate of detections on the polished surface that are not matched on the back surface.

Granite is often very old. A typical value for the exposure time might be 100 – 500 million years, although much older rocks are widely known. We adopt $T_e = 500$ million years as a fiducial value, but the precise region of parameter space that can be probed will depend on the provenance of the granite slabs.

In Figure 1, all three regions assume a minimum feature diameter of 1 mm. (We assume circular features for simplicity.) The first region with diagonal hatching assumed a slab area of just 10m^2 , the second, unatched region assumed a slab of area 100m^2 and The third region, with horizontal hatching, assumed a slab area of 1000m^2 .

In Figure 2, we have shown the range of parameter space that will be probed by examining a fixed area of granite but by placing various requirements on the sized of the tracks left by a passing macro. The region with diagonal hatching assumes searches for features with diameter of 0.1 cm or more, the unatched region is for a minimum feature diameter of 1 cm, and the horizontal-hatched region assumes a minimum feature diameter of 10 cm.

It is likely that the granite slabs we will use for this search will have somewhat different grain sizes and that the region probed will be a convolution of the various

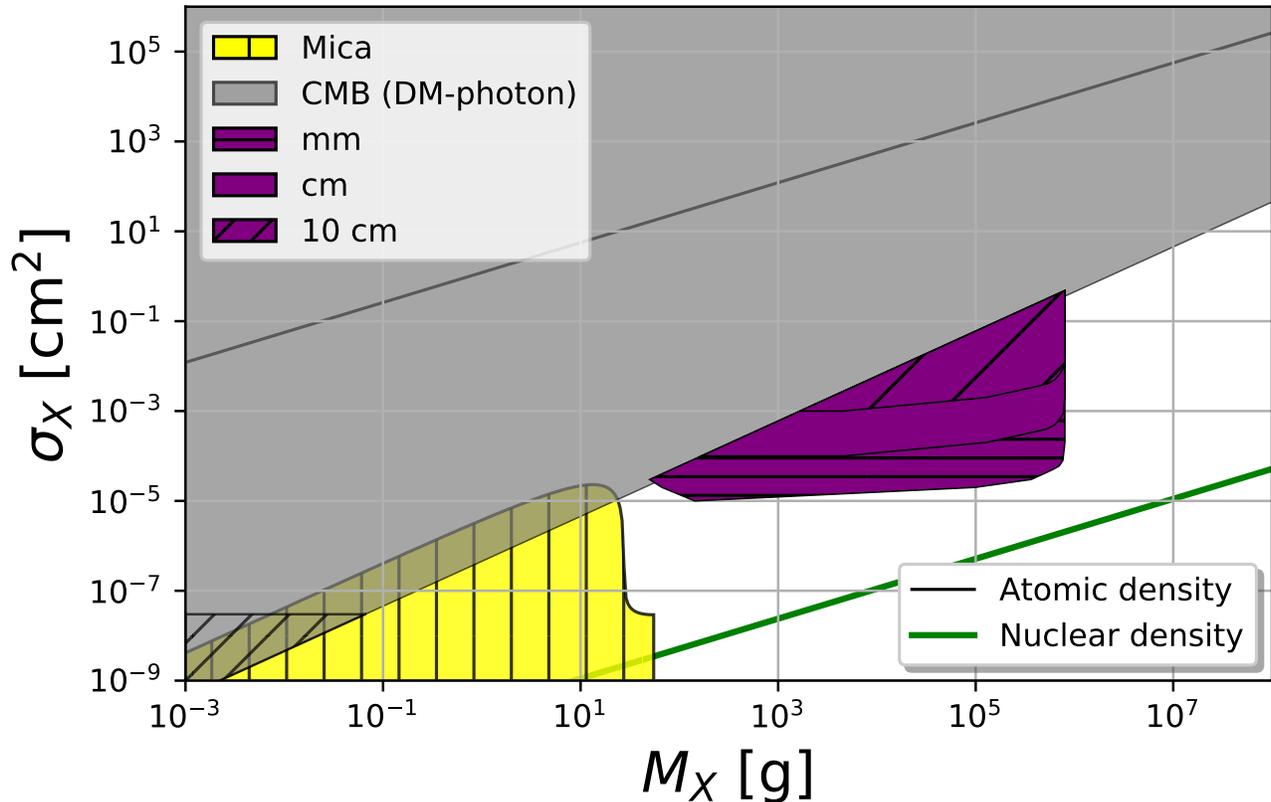


Figure 2. Illustrative regions of parameter space that could be probed from observations of slabs of ordinary rock are shown in purple. All 3 regions assume a slab area of 1000 m^2 . The region with diagonal hatching assumes searches for features with diameter of 0.1 cm or more, the unhatched region is for a minimum feature diameter of 1 cm, and the horizontal-hatched region assumes a minimum feature diameter of 10 cm. The region currently excluded by examination of ancient mica [11, 12] is shown in yellow with vertical hatching. The grey region is excluded from the effects of cosmic microwave background photons scattering off the macros [19]. The diagonal hatching of this bound in the bottom-left of the parameter space represents the region where the bound may weaken. Lines corresponding to nuclear and atomic density are shown for illustrative purposes.

regions corresponding to different sized features.

The calculation presented above assumed that the macro moves through the rock at a velocity equal to its impact velocity on the surface. At sufficiently high σ_X , this will no longer be true. In Figure 1, we insist that the σ_X be small enough for a macro that hits the Earth’s surface at 250 km s^{-1} to penetrate 5 km of rock without slowing to 100 km s^{-1} . If we were able to confidently identify rock samples that remained closer to the surface, or were older, we would be able to probe to larger values of σ_X and M_X . (However, we expect that at densities approaching atomic density, the macro, like a meteoroid, would disintegrate in the atmosphere or on impact with the Earth’s surface.) A full experimental analysis would integrate over the possible incidence angles of the macros and the velocity distribution of the dark matter.

We see that a manageable search for features that

can easily be identified by eye, in a quantity of granite slabs such as are normally found at a typical commercial countertop showroom, will begin to probe unexplored regions of parameter space, but not down to nuclear density. Since we do not know the detailed microphysics of macros, it is valuable to probe all open parameter values. Moreover, this search would serve as an important proof-of-concept for scaling up to the large-scale effort that would be required to push a search for macro dark matter down to the nuclear-density line.

ACKNOWLEDGMENTS

This work was partially supported by Department of Energy grant de-sc0009946 to the particle astrophysics theory group at CWRU.

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